

Potent strategy towards strongly emissive nitroaromatics through a weakly electron-deficient core†

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Nitroaromatics seldom fluoresce. The importance of electron-deficient (n-type) conjugates, however, has inspired a number of strategies for suppressing the emission-quenching effects of the strongly electron-withdrawing nitro group. Here, we demonstrate how such strategies yield fluorescent nitroaryl derivatives of dipyrromethane (DPND). Nitro groups near the DPND core quench its fluorescence. Conversely, nitro groups placed farther from the core allow some of the highest fluorescence quantum yields ever recorded for nitroaromatics. This strategy of preventing the known processes that compete with photoemission, however, leads to the emergence of unprecedented alternative mechanisms for fluorescence quenching, involving transitions to dark  $n\pi^*$  singlet states and aborted photochemistry. Forming  $n\pi^*$  triplet states from  $\pi\pi^*$  singlets is a classical pathway for fluorescence quenching. In nitro-DPNDs, however, these  $n\pi^*$  and  $\pi\pi^*$  excited states are both singlets, and they are common for nitroaryl conjugates. Understanding the excited-state dynamics of such nitroaromatics is crucial for designing strongly fluorescent electron-deficient conjugates.

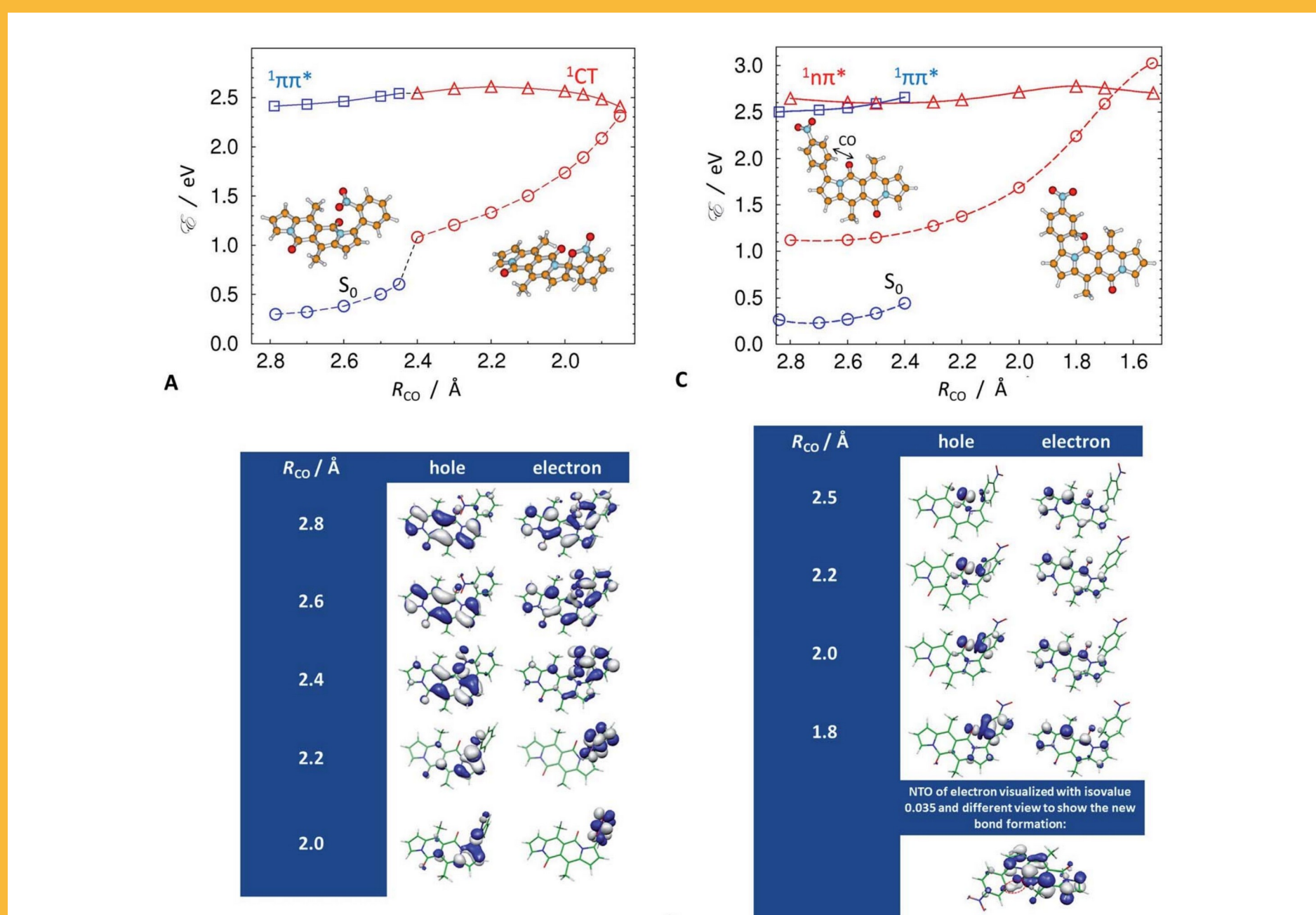


Fig. 9 Potential-energy profiles (PEPs) along the carbon-oxygen distances,  $R_{CO}$ , and the corresponding  $S_1 \rightarrow S_0$  NTOs (isovales = 0.03, unless indicated otherwise) optimized with the ADC(2)/def-SVP method in the lowest excited singlet states of (A) and (B) 5 and (C) and (D) 3. (A and C) Circles connected by dashed line denote vertical energy of the ground state computed at the geometry of the respective excited state. For 5, (A) locally-excited,  $\pi\pi^*$  state – blue squares, and the dark charge-transfer, CT, state – red triangles; and (B) NTOs along  $R_{CO}$ . For 3, (C) locally-excited,  $\pi\pi^*$  state – blue squares, and the dark  $n\pi^*$  state – red triangles; and (D) NTOs along  $R_{CO}$ , where the last bottom structure depicts the excited NTO at isovalue of 0.035 to visualize the formation of the carbon-oxygen covalent bond, as circled with a red dashed line.

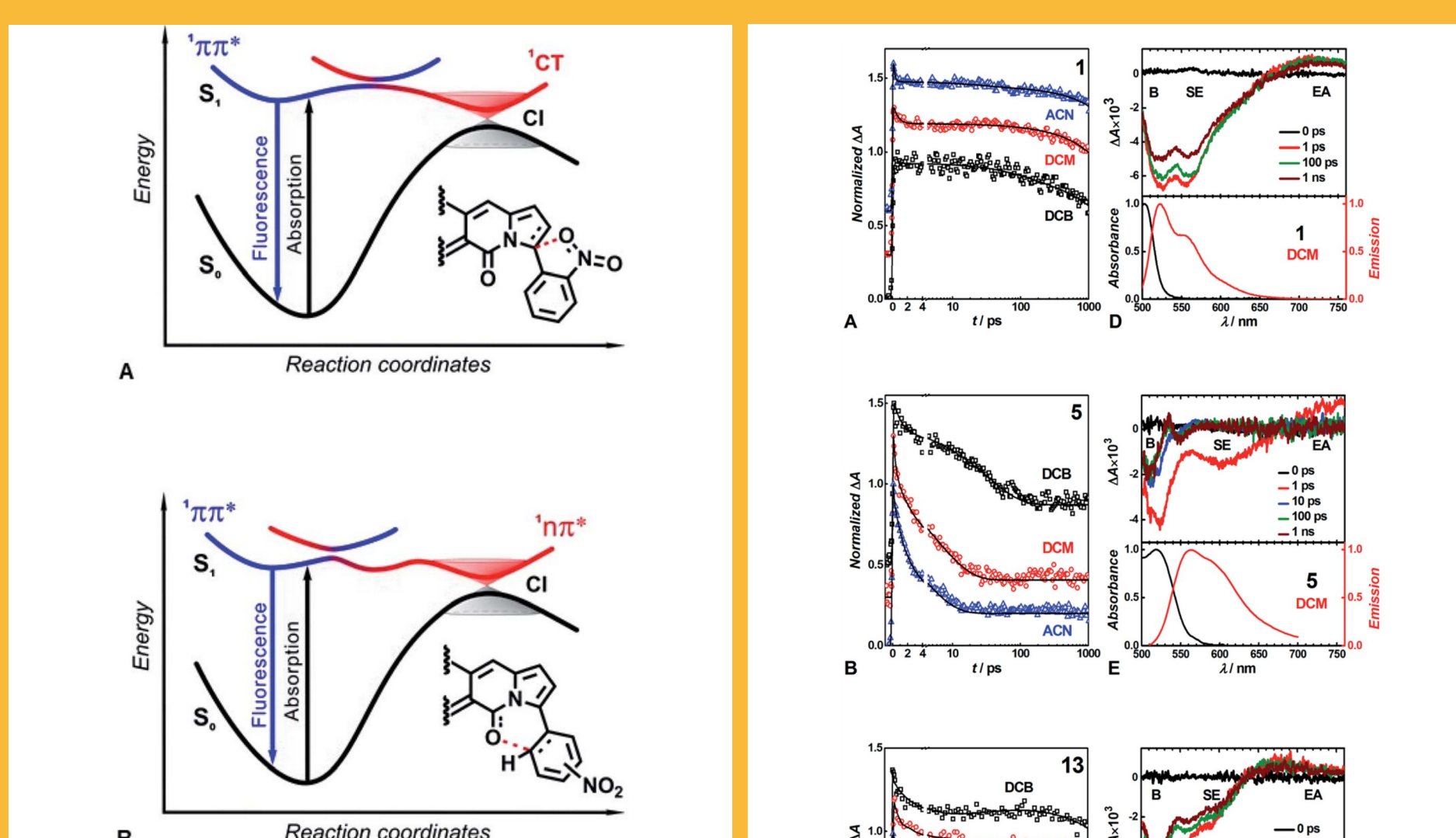


Fig. 10 Excited-state dynamics of nitroaryl-substituted DPNDs. The formed LE  $\pi\pi^*$  states either radiatively decay to  $S_0$ , or non-adiabatically transfer to dark states that form CIs with  $S_0$ , providing pathways for non-radiative deactivation. (A) Placing the nitro groups close to the DPND core, such as in the ortho and peri derivatives, accommodates the formation of dark CT states and the non-radiative deactivation assumes a  $\pi\pi^* \rightarrow {}^1CT \rightarrow S_0$  pathway. (B) When the nitro groups are not spatially close to the DPND core, such as in the para and meta derivatives, the energy level of the FC second singlet excited state, with a  $n\pi^*$  character, decreases along the reaction coordinate, opening a  $\pi\pi^* \rightarrow n\pi^* \rightarrow S_0$  pathway for non-radiative deactivation.

Fig. 6 Transient absorption decays normalized to 1 of (A) 1, (B) 5 and (C) 13 in different solvents at observation wavelength of 520 nm. Transient absorption spectra at different pump-probe time delays (upper panel) and time-integrated absorption and emission spectra normalized to 1 (lower panel) of (D) 1, (E) 5 and (F) 13 in DCM. The transient absorption signals were recorded upon excitation at 480 nm. Solid lines are from global multiplexponential fits to the experimental data. For clarity the transient absorption decays were offset on the y-axis. B – bleach band, SE – stimulated emission, EA – excited state absorption.



Combining Perovskites and Quantum Dots: Synthesis, Characterization, and Applications in Solar Cells, LEDs, and Photodetectors

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Metal halide perovskites having high defect tolerance, high absorption characteristics, and high carrier mobility demonstrate great promise as potential light harvesters in photovoltaics and optoelectronics and have experienced an unprecedented development since their occurrence in 2009. Semiconductor quantum dots (QDs), on the other hand, have also been proved to be very flexible toward shape, dimension, bandgap, and optical properties for constructing optoelectronic devices. Of late, a strategic combination of both materials has demonstrated extraordinary promise in photovoltaic applications and optoelectronic devices. Combining QDs and perovskites has proved to be quite an effective strategy toward the formation of pinhole-free and more stable perovskite crystals along with tunability of other properties. To boost this exciting research field, it is imperative to summarize the work done so far in recent years to provide an intriguing insight. This review is a critical account of the advanced strategy toward combining these two fascinating materials, including their different synthetic approaches regarding heteroepitaxial growth of perovskite crystals on QDs, carrier dynamics at the interface and potential application in the fields of solar cells, light emitting diodes, and photodetectors.

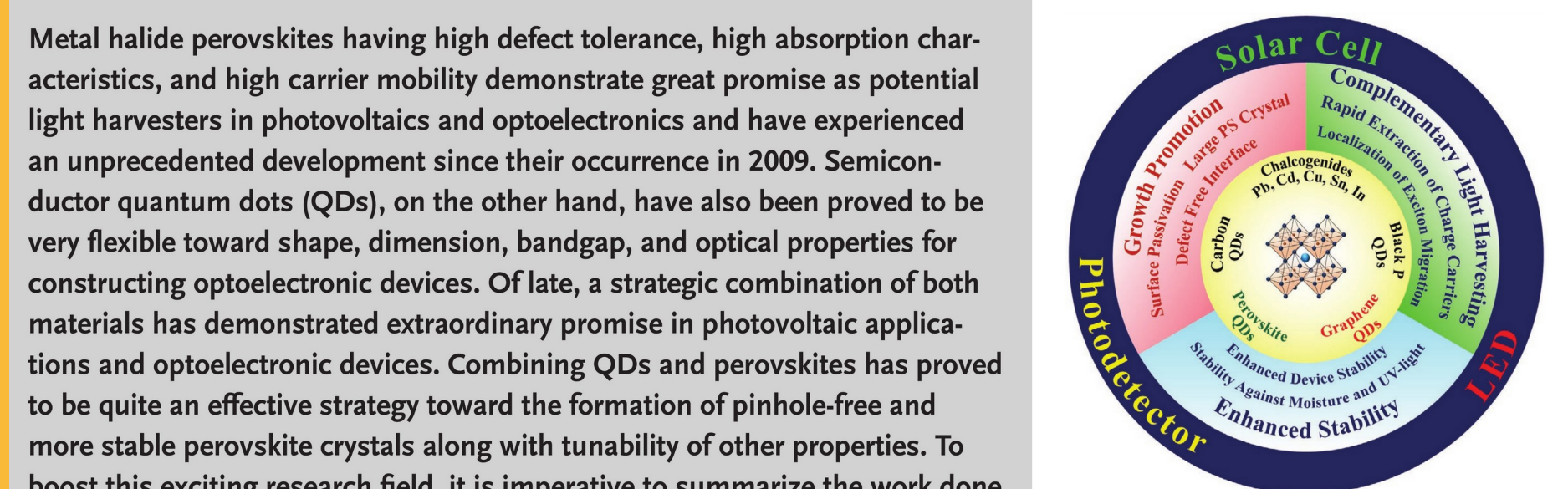


Fig. 1 Schematic illustration of the effect of different integrated quantum dots on different properties of perovskites and their applications in the QD/perovskite hybrid systems.



Effect of the Mixture Composition of BmimBF4-Acetonitrile on the Excited-State Relaxation Dynamics of a Solar-Cell Dye D149: An Ultrafast Transient Absorption Study

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ABSTRACT: The excited-state relaxation dynamics of D149, one of the metal-free substituted indoline dyes used in dye-sensitized solar cells, is studied in the whole composition range of the 1-butyl-3-methylimidazolium tetrafluoroborate (BmimBF<sub>4</sub>)-acetonitrile binary mixture by using time-integrated absorption, emission, and time-resolved transient absorption (TA) spectroscopies. The comparative analysis of absorption and emission spectra indicates that the value of Stokes shift reduces monotonically with decreasing mixture polarity. The global analysis of time-resolved TA spectra indicates the presence of four different time components related to different processes in the excited state of the dye. Importantly, the observed timescales are highly sensitive to composition, polarity, and viscosity of the binary mixture. Increase of viscosity and decrease of polarity observed for increasing ionic liquid (IL) content in the mixture lead to overall increase in the emission lifetime ( $S_1 \rightarrow S_0$ ) of D149. At a lower IL mole fraction ( $X_{IL} = 0.1$ ), the emission lifetime shows a minimum that can be traced back to the change from the situation in which the local environment of the dye is dominated by the interactions in acetonitrile to that in which it is dominated by those in BmimBF<sub>4</sub>. This also is reflected in the occurrence of a minimum in relative quantum yield in the same range of  $X_{IL}$ . The origin of the other moderately long-time component (33 ps in ACN/120 ps in BmimBF<sub>4</sub>) is still debatable; however, for pure IL and all the mixtures, the composition dependence of this timescale is similar to that of the longest emission lifetime.

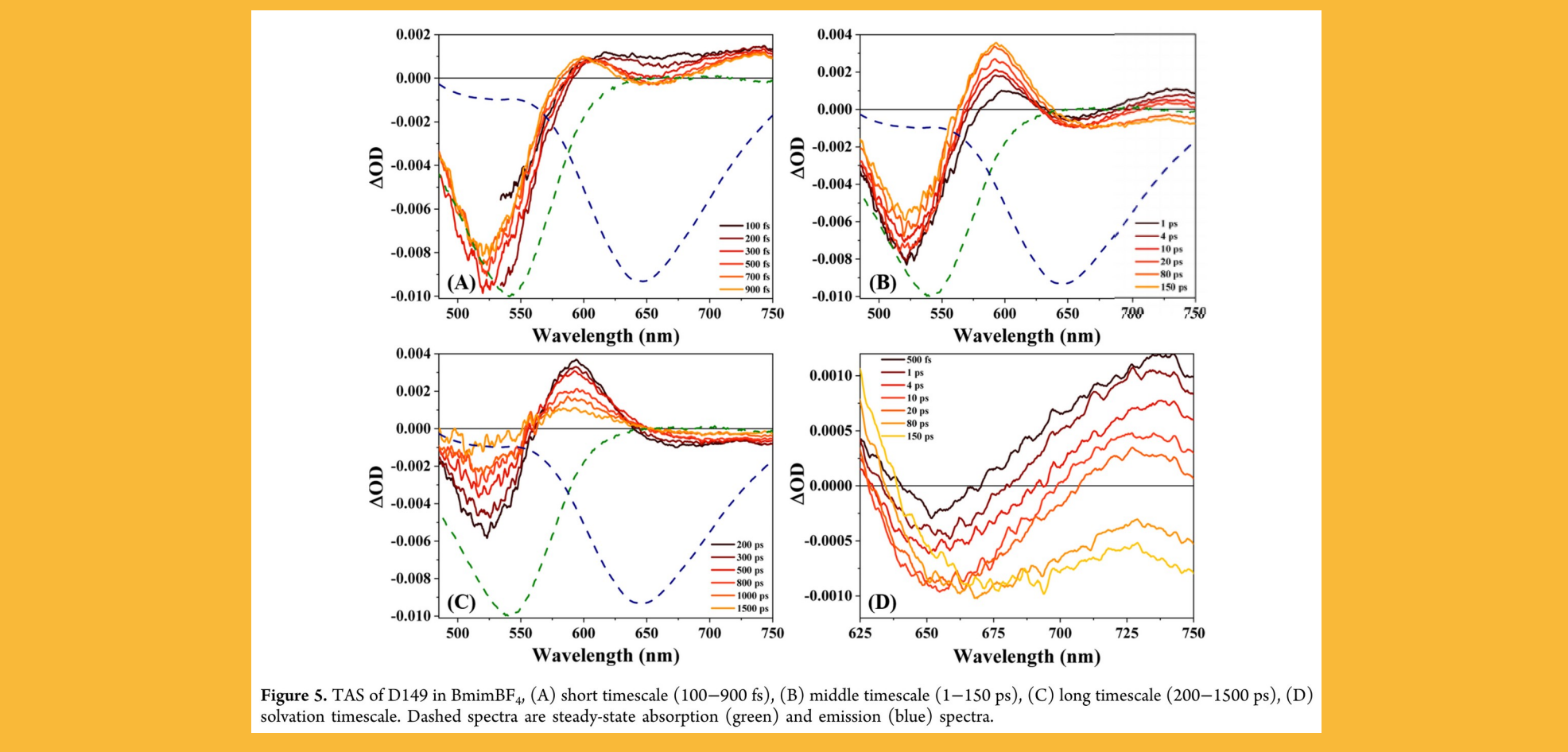


Figure 8. TAs of D149 in BmimBF<sub>4</sub> (A) short timescale (100–900 fs), (B) middle timescale (1–150 ps), (C) long timescale (200–1500 ps), (D) solvation timescale. Dashed spectra are steady-state absorption (green) and emission (blue) spectra.

Dynamics in the BMIM PF<sub>6</sub>/acetonitrile mixtures observed by femtosecond optical Kerr effect and molecular dynamics simulations†

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We have performed the measurements of the optical Kerr effect signal time evolution up to 4 ns for a mixture of 1-alkyl-3-methylimidazolium hexafluorophosphate (BMIM PF<sub>6</sub>) ionic liquid and acetonitrile in the whole mole fractions range. The long delay line in our experimental setup allowed us to capture the complete reorientational dynamics of the ionic liquid. We have analyzed the optical Kerr effect signal in the time and frequency domains with help of molecular dynamics simulations. In our approximation of the slow picosecond dynamics with a multi-exponential decay, we distinguish three relaxation times. The highest two are assigned to the reorientation of the cation and acetonitrile molecules that are in the vicinity of the imidazolium ring. The third one is recognized as originating from cation rotations and reorientation of acetonitrile molecules in the bulk or in the vicinity of the aliphatic chains of the cation. With help of the simulation we interpret the intermolecular band in the reduced spectral density, obtained from Kerr signal, as follows: its low-frequency side results from oscillations of one of the components in the cage formed by its neighbors, while the high-frequency side is attributed to the librations of the cation and acetonitrile molecule as well as the intermolecular oscillations of system components involved in specific interactions. We use this assignment and concentration dependence of the spectra obtained from velocity and angular velocity correlations to explain the mole fraction dependence of Kerr reduced spectral density.

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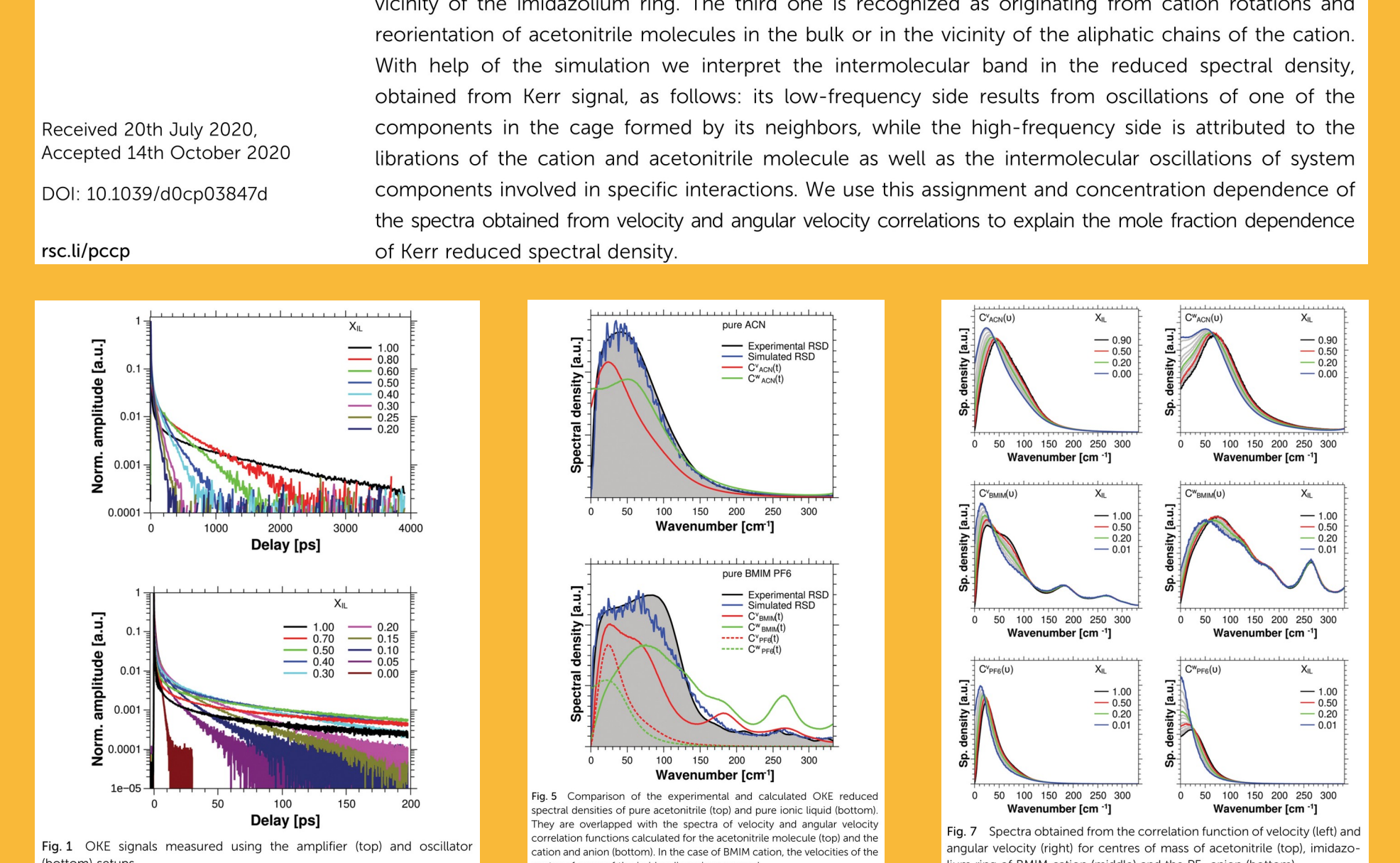


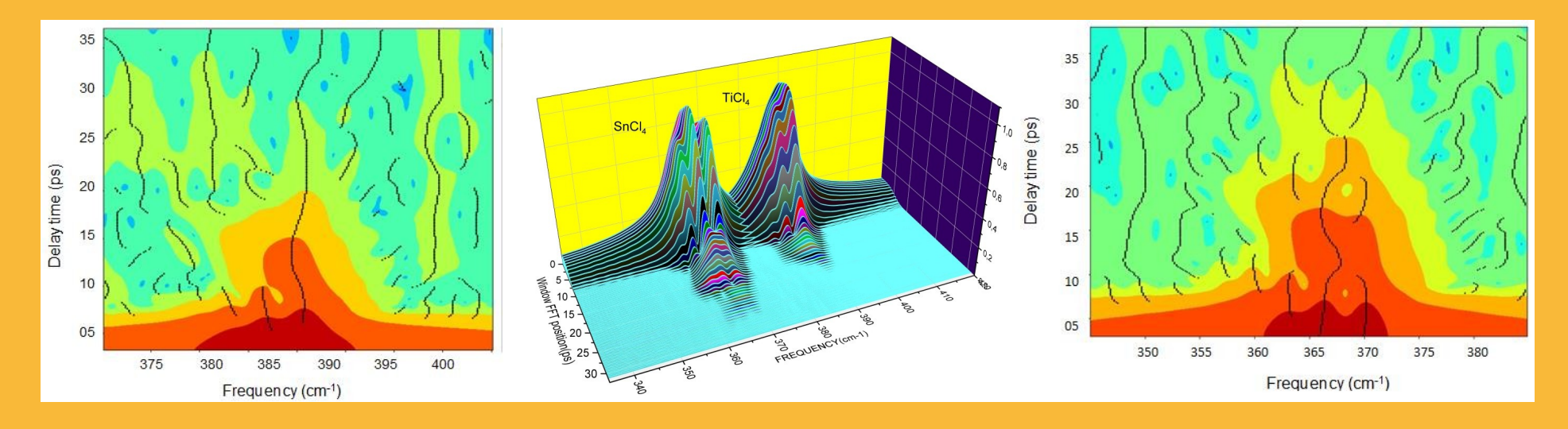
Fig. 9 Comparison of the experimental and calculated optical Kerr effect signal. The experimental signal is shown in black, the calculated signal is shown in red. The inset shows the calculated signal for different mole fractions of BMIM PF<sub>6</sub> in acetonitrile.

Time resolved transient transmission spectroscopy of TiCl<sub>4</sub> and SnCl<sub>4</sub>

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Abstract  
Herewith, for the first time, we present the vibrational spectra collected for liquid TiCl<sub>4</sub> and liquid SnCl<sub>4</sub> by use of time resolved transient transmission spectroscopy. Of our interest is the isotopically split isotropic intramolecular vibrational band, the shape of which is very sensitive to intermolecular interactions. The high resolution spectra, obtained as fast Fourier transforms of the time domain signals acquired in transient transmission experiment, are compared with spontaneous Raman spectra. The dependence of the spectrum shape on intermolecular interactions has been established by diluting TiCl<sub>4</sub> and SnCl<sub>4</sub> in CS<sub>2</sub> at different concentrations. Fitting the simplified oscillator model of a liquid to FFTs of time domain signals in transient transmission experiment we have found intermolecular force constants for all concentrations. Application of the pump-probe spectroscopy technique and windowed fast Fourier transform procedure allowed us to observe the evolution of the spectral shape, and thus of the intermolecular forces, after the liquid has been perturbed by the femtosecond pump pulse.



Search for the origin of synergistic solvation in methanol/chloroform mixture using optical Kerr effect spectroscopy

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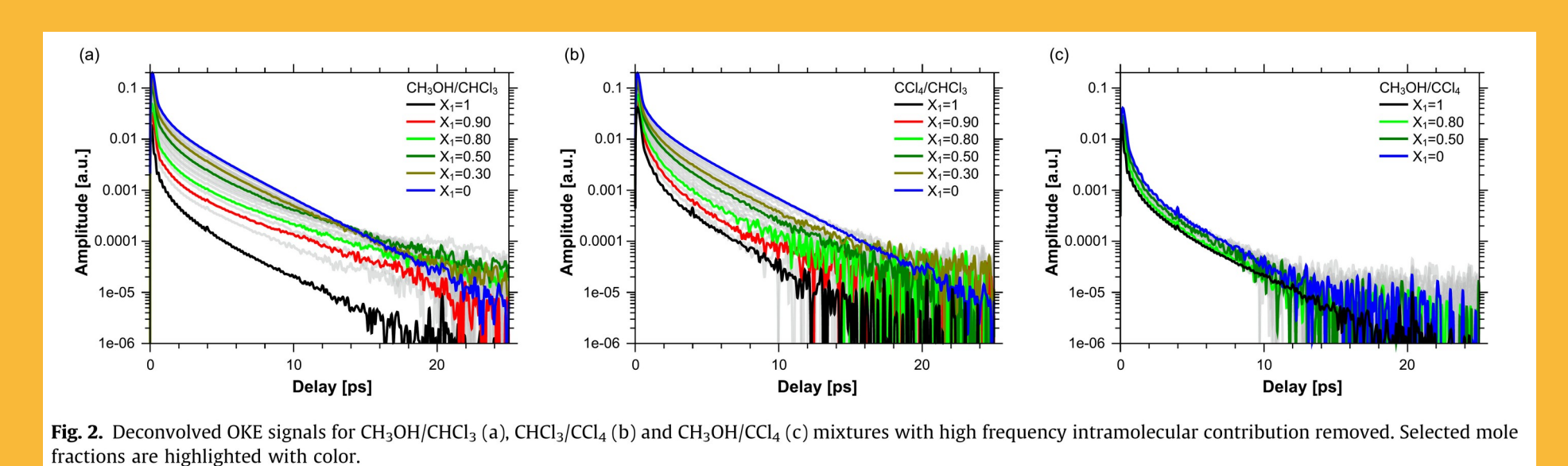


Fig. 2. Deconvoluted OKE signals for CH<sub>3</sub>OH/CHCl<sub>3</sub> (a), CHCl<sub>3</sub>/CCl<sub>4</sub> (b) and CH<sub>3</sub>OH/CCl<sub>4</sub> (c) mixtures with high frequency intramolecular contribution removed. Selected mole fractions are highlighted with color.

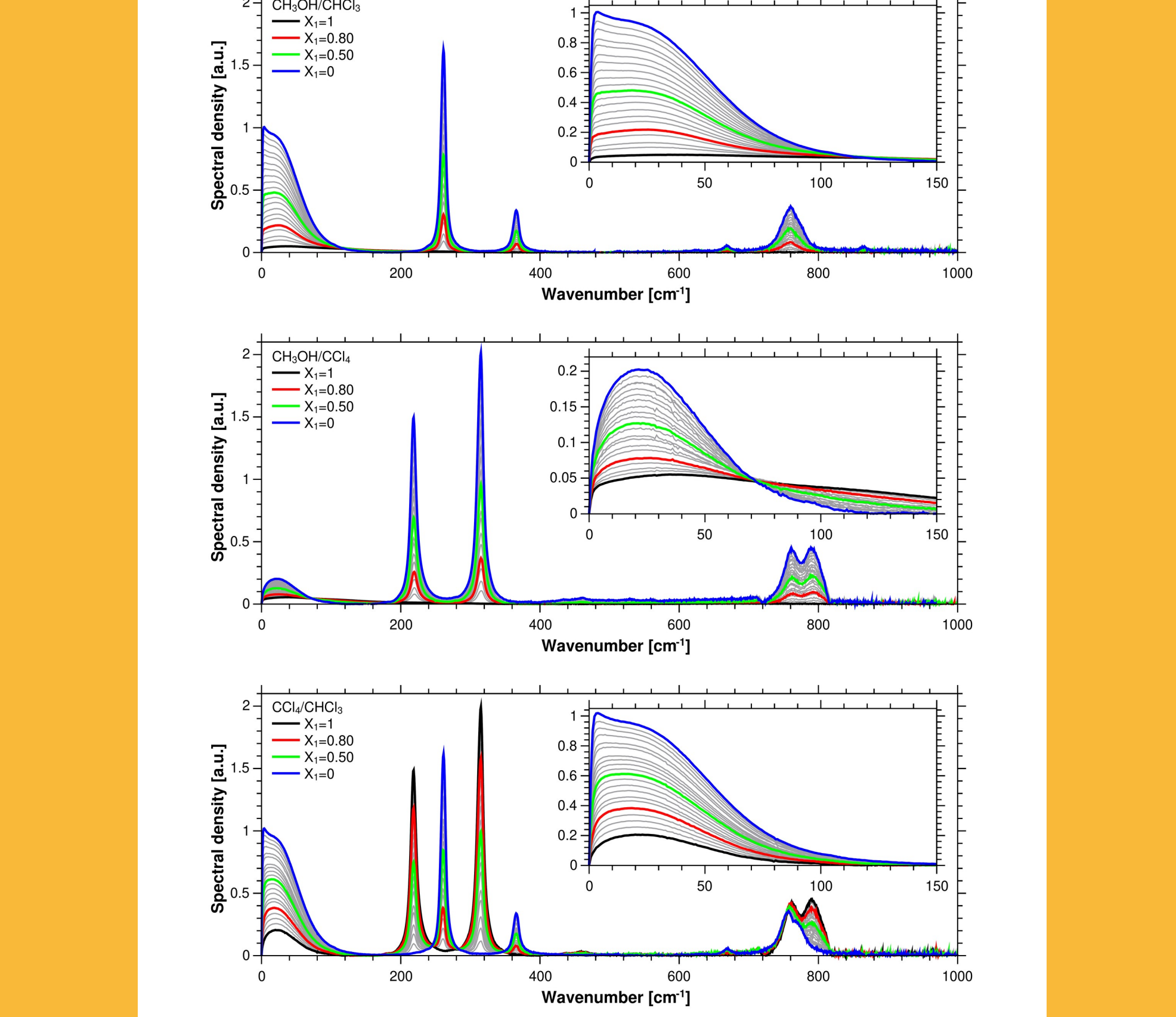


Figure S22. Mole fraction dependence of the OKE total spectral density for CH<sub>3</sub>OH/CHCl<sub>3</sub>, CH<sub>3</sub>OH/CCl<sub>4</sub> and CCl<sub>4</sub>/CHCl<sub>3</sub> mixtures. The insets show the intermolecular band in the 0–150 cm<sup>-1</sup> range.

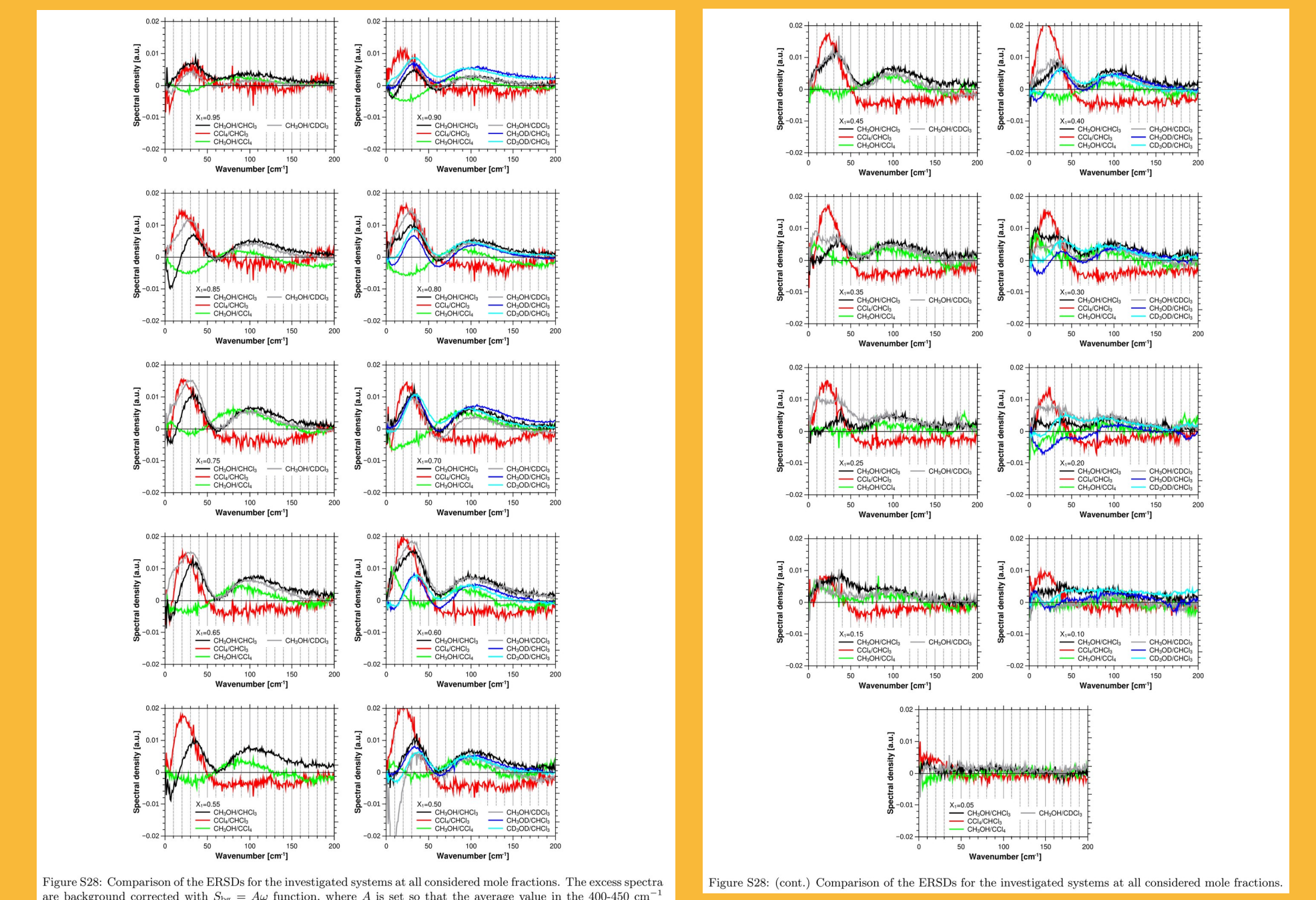


Figure S20. Comparison of the ESDs for the investigated systems at all considered mole fractions. The axes are the same as in Figure S22. The insets show the intermolecular band in the 0–150 cm<sup>-1</sup> range.

